SEPARATION THERMODYNAMICS

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Summary

In this concise article, the relevant thermodynamic functions are defined for the separation of components from their mixtures. Seawater or brine being a mixture of salts and water, the equilibrium criteria are important for consideration in the desalination operations.

1. Relevant Thermodynamic Functions

The following is the fundamental equation interrelating primary thermodynamic variables for a single-phase system from which all other equations connecting the properties of such systems are derived:

$$d(nU) = Td(nS) - Pd(nV) + \Sigma(\mu_i dn_i)$$
(1)

The quantity μ_i , which is the chemical potential of component i, plays a vital role in the thermodynamics of phase and chemical equilibria. Equation (1) suggests possible combinations of terms that may be defined as primary functions; the most important are the following.

Internal energy: U

Enthalpy:
$$H = U + PV$$
 (2)

Helmholtz function:
$$A = U - TS$$
 (3)

Gibbs function (also called free energy):
$$G = U + PV - TS = H - TS$$
 (4)

For one mole of a homogeneous fluid having a constant composition, the following relations can be written

$$dU = TdS - PdV (5)$$

$$dH = TdS + VdP \tag{6}$$

$$dA = -PdV - SdT \tag{7}$$

$$dG = VdP - SdT \tag{8}$$

Next, fugacity is defined in direct relation to the Gibbs function; for a constant composition mixture

$$dG = RTd \ln f \quad \text{(constant } T, x) \tag{9}$$

with
$$\lim_{P \to 0} \frac{f}{P} = 1$$
.

It can be shown that when the ideal gas law holds, fugacity becomes equal to pressure.

2. Partial Molar Properties

A homogeneous system may contain any number of components or chemical species. If M represents a molar thermodynamic property of the mixture or solution, such as U, H, S, etc., the total system property becomes nM, where

$$n = \sum_{i} n_{i}$$

are the total moles in the system. It may be expected to relate the mixture property M to the properties M_i of the pure chemical species making up the mixture. However, no generally valid rules are known and the connection must be established experimentally

for a particular system.

The total mixture property nM is extensive, hence, it may be convenient to apportion this property among the mixture constituents in an appropriate way. It follows from Euler's theorem on homogeneous functions that

$$nM = \sum \left[n_i \left(\frac{\partial (nM)}{\partial n_i} \right)_{T,P,n_j} \right]$$
 (10)

from which a partial molar property

 \overline{M}_i , of component i in the solution can be defined as follows

$$\bar{M}_{i} = \left[\frac{\partial(nM)}{\partial n_{i}}\right]_{T,P,n_{i}} \tag{11}$$

Substituting Eq. (11) into Eq. (10),

$$nM = \Sigma(n_i \overline{M}_i) \tag{12}$$

The \overline{M}_i s, are homogeneous functions of zero degree in the mole numbers; therefore, they are intensive thermodynamic properties. Thus, Eq. (11) determines how the mixture property can be apportioned among the mixture constituents and Eq. (12) shows that the total mixture property is the sum of component properties. Dividing Eq. (12) by n converts it on a mole basis, i.e.

$$M = \sum (x_i \overline{M}_i) \tag{13}$$

where x_i represents the mole fraction of i in the mixture.

The equations for partial properties can be applied on a unit mass basis as well. In this case, mass m replaces moles n and x_i s represent mass fraction. Furthermore, it is readily shown that

$$\overline{H}_i = \overline{U}_i + P\overline{V}_i; \overline{A}_i = \overline{U}_i - T\overline{S}_i; \overline{G}_i = \overline{H}_i - T\overline{S}_i.$$

Equation (12) is a general relationship for any homogeneous system in an equilibrium state. Therefore, changes in nM caused by changes in P, T, or the n_i s are given by the total differential

$$d(nM) = \left[\frac{\partial(nM)}{\partial T}\right]_{P,n} dT + \left[\frac{\partial(nM)}{\partial P}\right]_{T,n} dP + \sum(\overline{M}_i \partial n_i)$$
(14)

Furthermore,

$$d(nM) = \sum_{i} (n_i d\overline{M}_i) + \sum_{i} (\overline{M}_i dn_i)$$
(15)

From Eqs (14) and (15),

$$n\left(\frac{\partial M}{\partial T}\right)_{P,x} dT + n\left(\frac{\partial M}{\partial P}\right)_{T,x} dP - \sum_{i} \left(n_{i} d\overline{M}_{i}\right) = 0$$
(16)

Dividing Eq. (16) by n,

$$\left(\frac{\partial M}{\partial T}\right)_{P,x} dT + \left(\frac{\partial M}{\partial P}\right)_{T,x} dP - \sum (x_i d\overline{M}_i) = 0$$
(17)

Equation (17) is known as the Gibbs-Duhem equation in the general form and is valid for any thermodynamic property M in a homogeneous phase. At constant T and P, it becomes

$$\sum (x_i d\overline{M}_i) = 0 \ (P \ \text{and} \ T \ \text{constant})$$
 (18)

The chemical potential μ_i is expressed as

$$\mu_i = \left| \frac{\partial (nG)}{\partial n_i} \right| \tag{19}$$

Comparing Eq. (19) with Eq. (11),

$$\mu_i = \overline{G}_i$$
.

Thus, the chemical potential is identical to the partial molar Gibbs function. In the pattern of Eq. (1) for other thermodynamic properties such as the Gibbs function and substituting

$$\mu_i = \overline{G}_i$$

it can be shown that

$$-SdT + VdP - \sum_{i} (x_i d\overline{G}_i) = 0$$
(20)

Equation (20) is the most commonly encountered form of the Gibbs-Duhem equation.

2.1. Excess Properties

An excess property of solution is defined as the difference between the actual value of

the property and the one that is calculated at the same condition of temperature, pressure, and composition for an ideal solution. Thus,

$$M^E = M - M_{\text{ideal}} \tag{21}$$

and

$$\Delta M^E = \Delta M - \Delta M_{\text{ideal}} \tag{22}$$

where M^E is the excess solution property and ΔM^E is the excess property change of mixing. Similarly, there are partial molar excess properties

$$\bar{M}_{i}^{E} = \bar{M}_{i} - \bar{M}_{i \text{ ideal}} \tag{23}$$

$$\bar{\Delta}\bar{M}_{i}^{E} = \bar{\Delta}\bar{M}_{i} - \bar{\Delta}\bar{M}_{i,\text{ideal}} \tag{24}$$

It can be shown that M^E and ΔM^E are identical, as are

$$\overline{M}_i^E$$
 and $\overline{\Delta}\overline{M}_i^E$.

If M is a molar thermodynamic property of a homogeneous fluid mixture, then ΔM is defined by

$$\Delta M = M - \sum (x_i M_i^o), \tag{25}$$

where ΔM is called property change in mixing and

$$M_i^o$$

is the molar property of the pure component i at the temperature of the mixture and in some standard state of specified pressure and composition. For convenience, the standard state is chosen as that of pure i in its stable state at the mixture pressure. Substituting from Eq. (13) into Eq. (25),

$$\Delta M = \sum (x_i \overline{M}_i) - \sum (x_i M_i^o) = \sum [x_i (\overline{M}_i - M_i^o)]$$

or

$$\Delta M = \sum_{i} (x_i \overline{\Delta M}_i), \tag{26}$$

where, by definition,

$$\overline{\Delta}\overline{M}_i \equiv \overline{M}_i - M_i^O$$
.

For species i in solution, the fugacity

 \bar{f}_i

is defined as

$$d\overline{G}_i = RTd \ln \overline{f}_i \quad \text{(constant } T), \tag{27}$$

with
$$\lim_{P \to 0} \bar{f}_i / (x_i P) = 1$$

If Eq. (27) is integrated for the change of state from pure i in a standard state to a constituent in solution at the same T, then

$$\overline{G}_i - G_i^o = RT \ln \frac{\overline{f}_i}{f_i^o} \tag{28}$$

and

$$\overline{f_{i \text{ ideal}}} = x_i \ f_i^o \tag{29}$$

From the definition of

 $\overline{\Delta}\overline{M}_i$

and Eq. (29), Eq. (28) becomes

$$\overline{\Delta}\overline{G}_{i,\text{ideal}} = RT \ln x_i \tag{30}$$

and from Eqs (26) and (30)

$$\Delta G_{\text{ideal}} = RT \sum x_i \ln x_i \tag{31}$$

and from Eqs (25) and (31)

$$G_{\text{ideal}} = \sum (x_i G_i^o) + RT \sum (x_i \ln x_i). \tag{32}$$

The most useful of these properties is the excess Gibbs function. Combining Eqs (22) and (31),

$$G^{E} = \Delta G - \Delta G_{\text{ideal}} = \Delta G - RT \sum_{i} (x_{i} \ln x_{i})$$
(33)

or

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$$G^{E} = G - \sum (x_i G_i^{\circ}) - RT \sum (x_i \ln x_i)$$
(34)

Dividing Eq. (34) by RT puts it into a dimensionless form:

$$\frac{G^{E}}{(RT)} = \frac{G}{(RT)} - \sum \left[\frac{x_i G_i^{o}}{(RT)} \right] - \sum (x_i \ln x_i). \tag{35}$$

2.2. Activity Coefficient

An auxiliary thermodynamic function is defined as the activity coefficient for a component in the solution, as follows:

$$\gamma_i = \frac{\overline{f_i}}{(x_i f_i^o)} \tag{36}$$

where

$$f_i^o$$

is the fugacity in the standard state. Writing Eq. (36) in the following form:

$$\ln \gamma_i = \ln \left(\frac{\overline{f_i}}{f_i^o} \right) - \ln x_i \tag{37}$$

Substituting from Eqs (28) and (30),

$$\ln \gamma_i = \frac{\overline{G}_i - G_i^o}{RT} - \frac{\overline{G}_{i,ideal} - G_i^o}{RT} = \frac{\overline{G}_i - \overline{G}_{i,ideal}}{RT}$$

or

$$\ln \gamma_i = \frac{\overline{G}_i^E}{(RT)} \tag{38}$$

Equation (38) shows that $\ln \gamma_i$ in relation to $G^E/(RT)$ stands as a partial molar property.

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